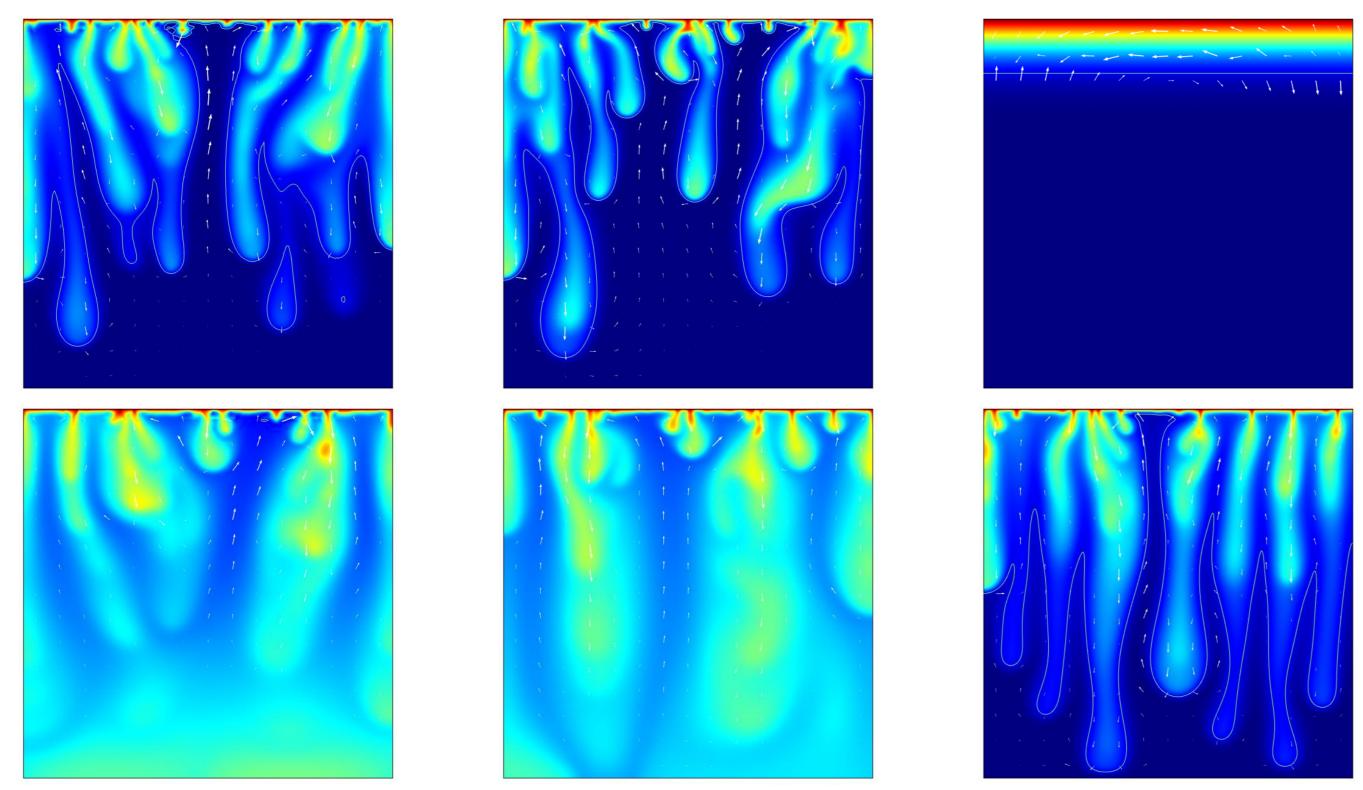
## Numerical Modelling of CO<sub>2</sub>-Storage

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**Introduction**: Storage of CO<sub>2</sub> in the sub-surface is seen as a technology that can contribute to the generally accepted goal of a de-carbonized society. As real field experiments are hardly feasible, many current studies utilize the capabilities of numerical modelling.

Concerning the practical application of  $CO_2$  storage many questions are still unanswered. In the most favoured scenario  $CO_2$  in supercritical state is pressed into a deep geological formation. Within the permeable layer  $CO_2$  will come to overlie brine and will start to dissolve into the deeper part by diffusion, which is influenced by convection.



**Figure 1**. CO<sub>2</sub> concentration distributions (red: saturated, blue: zero; top: early convection, bottom: late convection; left: coarse mesh, center: medium mesh, right: fine mesh)

Convection is a multi-physics phenomenon, in which flow and transport processes are coupled. For the coupling the fluid density is the crucial parameter. For the highly dynamic processes of  $CO_2$  storage, with high Rayleigh number, the initial phase with pure diffusion is followed by a convection phase. The latter can be sub-divided in an early stage with high but fluctuating mass transfer; and a late stage, in which mass transfer is decreasing<sup>2</sup>.

Computational Methods: Flow and transport are described by a non-linear set of two partial differential equations: (1) one for the streamfunction  $\Psi$ , which includes the dimensionless Rayleigh number Ra, and (2) for the normalized concentration  $c^{1,3}$ :

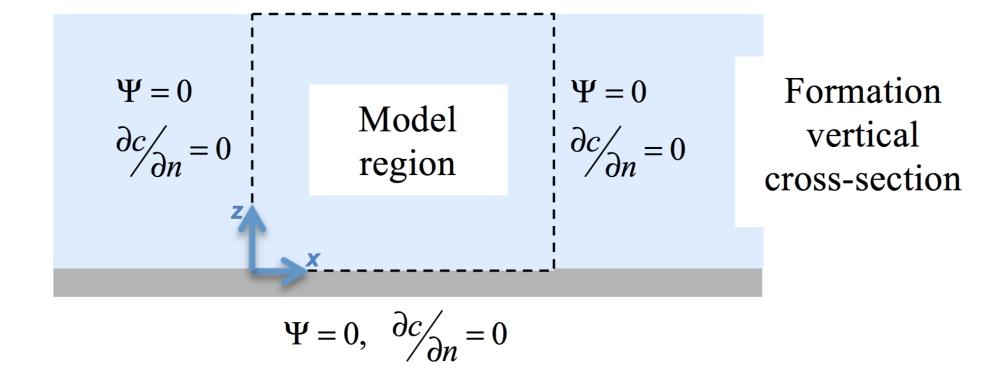
$$\frac{f}{fx}\left(\frac{f\Psi}{fx}\right) + \frac{f}{fz}\left(\frac{f\Psi}{fz}\right) = -Ra\frac{fc}{fx} \quad \text{with} \quad Ra = \frac{gk\Delta\rho H}{\mu D} \tag{1}$$

$$\frac{fc}{ft} = \nabla \cdot (\nabla c - \mathbf{v}c) \quad \text{with} \quad \mathbf{v}_{x} = -\frac{f\Psi}{fz} \quad \text{and} \quad \mathbf{v}_{z} = \frac{f\Psi}{fx}$$
 (2)

The following sketch shows the 2D model region within a vertical cross-section through the formation and the boundary conditions:  $\Psi=0,\ c=1$ 

Figure 2. Model region & boundary conditions



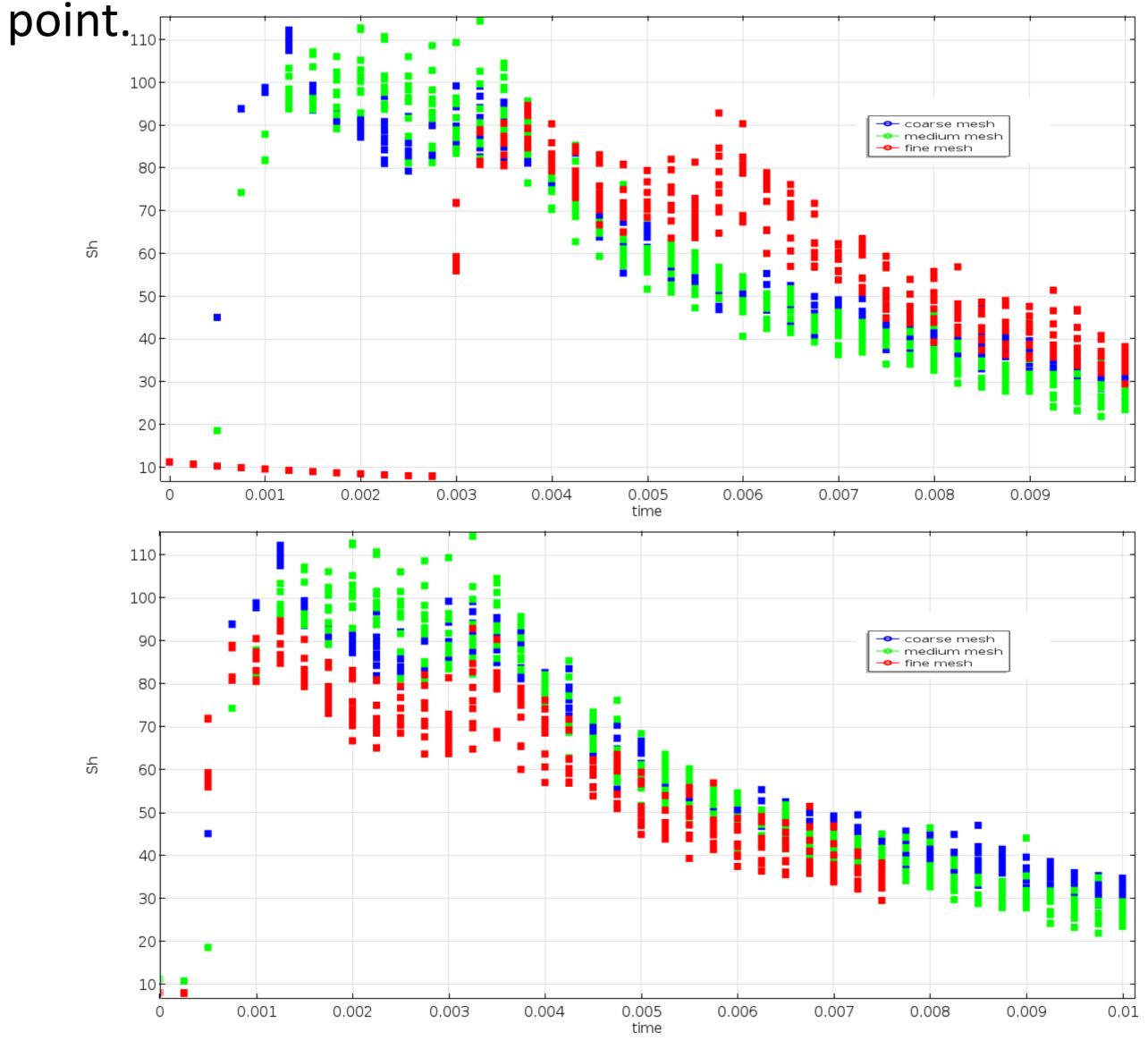


**Results**: Using COMSOL Multiphysics® the inhomogeneity of the permeability field, realized by random distributions were explored, as well as mesh dependencies. Input parameters are given in Table 1.

Parameter	Value [Unit]	Parameter	Value [Unit]
Saturated CO <sub>2</sub>	0.0493	Density difference	$10.45 \text{ kg/m}^3$
mass fraction		$\Delta  ho$	
Viscosity $\mu$	$0.5947 \cdot 10^{-3}$	Molecular	$2 \cdot 10^{-9} \text{ m}^2/\text{s}$
	Pa•s	diffusivity D	
Brine density	994.56 kg/m <sup>3</sup>	Permeability <i>k</i> (mean value)	5•10 <sup>-13</sup> m <sup>2</sup>

**Table 1**. List of reference case parameters<sup>4</sup>

The input Rayleigh number Ra was 5000. Fig. 3 shows the temporal development of mean mass transfer obtained from 10 realisations for 3 different meshes. The lower figure is obtained by shifting the fine mesh results in order to match the 'onset on convection'



**Figure 3**. Mass transfer range, indicated by the Sherwood number Sh for 30 different scenarios

Conclusions: The details of the flow patterns depend heavily on disturbances of physical parameters and on numerical features. Comparing mass transfer of 30 scenarios, a range for mass transfer during the different stages is obtained. Despite the differences in the single scenarios the duration of the early convection stage turns out to be a constant. Moreover, in the late convection stage the range of fluctuations is decreasing.

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