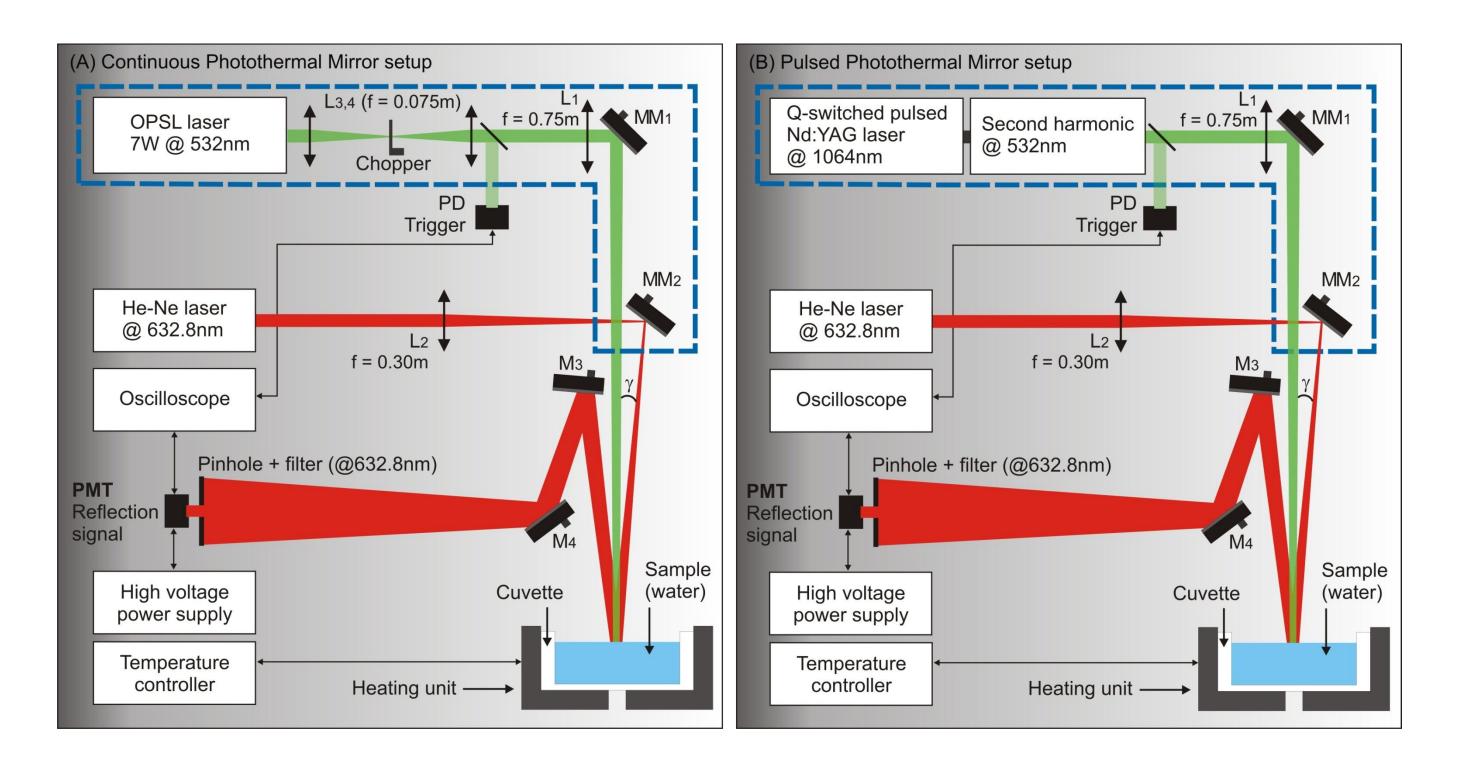
Radiation Force Effect at the Dielectric Water-Air Interface Gustavo V. B. Lukasievicz¹, Nelson G. C. Astrath², Luis C. Malacarne², Mauro L. Baesso², Stephen E. Bialkowski³ 1. Universidade Tecnológica Federal do Paraná, COINF, Toledo, Paraná, Brazil; 2. Universidade Estadual de Maringá, Departamento de Física, Maringá, Paraná, Brazil; 3. Utah State University, Department of Chemistry and Biochemistry, Logan, Utah, USA.

Introduction: The effects of radiation pressure exerted on a dielectric surface exposed to electromagnetic radiation has been a long-standing debate for over a century. The effect can be interpreted as the transfer of momentum from photons at the surface in the direction of propagation of the incident electromagnetic radiation^{1,2}.



Results: Figures 3 and 4 display the actual deformation of water at different exposure times¹. Under continuous (cw) excitation, the liquid surface rises with time reaching a maximum deformation of around 30nm at the center of the excitation beam, Figure 3. As for the pulsed excitation, a sharp peak appears at short time, which is dispersed rapidly on the surface, Figure 4.

Figure 1. Schematic diagram of the time-resolved photomechanical mirror under continuous (a) and pulsed (b) experiments¹.

Computational Methods: The radiation force effects on the surface displacement can be calculated by solving the Navier-Stokes equation with appropriated boundary conditions. The surface deformation can be described by the radiation pressure as well as those forces due to gravity and surface tension.

 $\mathbf{F} = -\rho g \hat{z}$

ρ

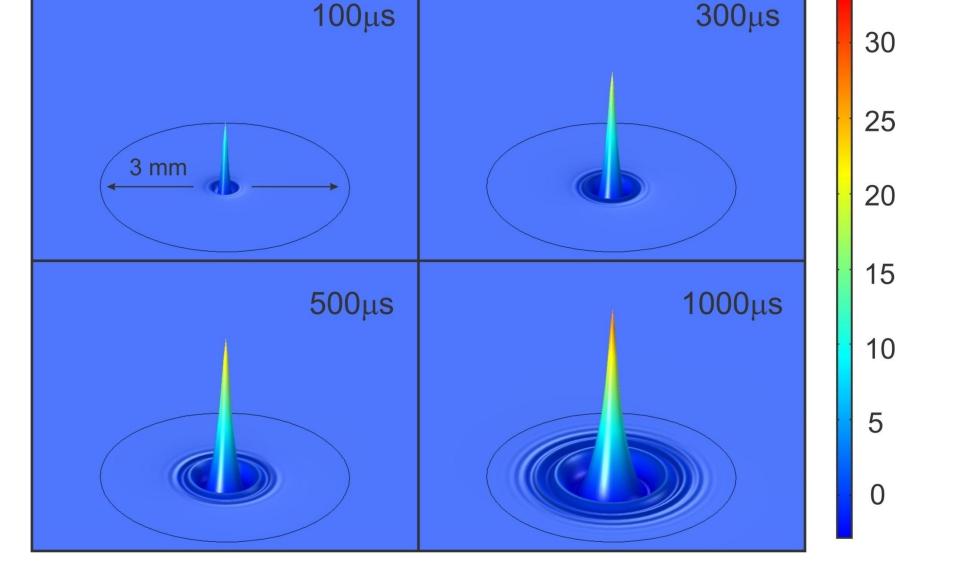
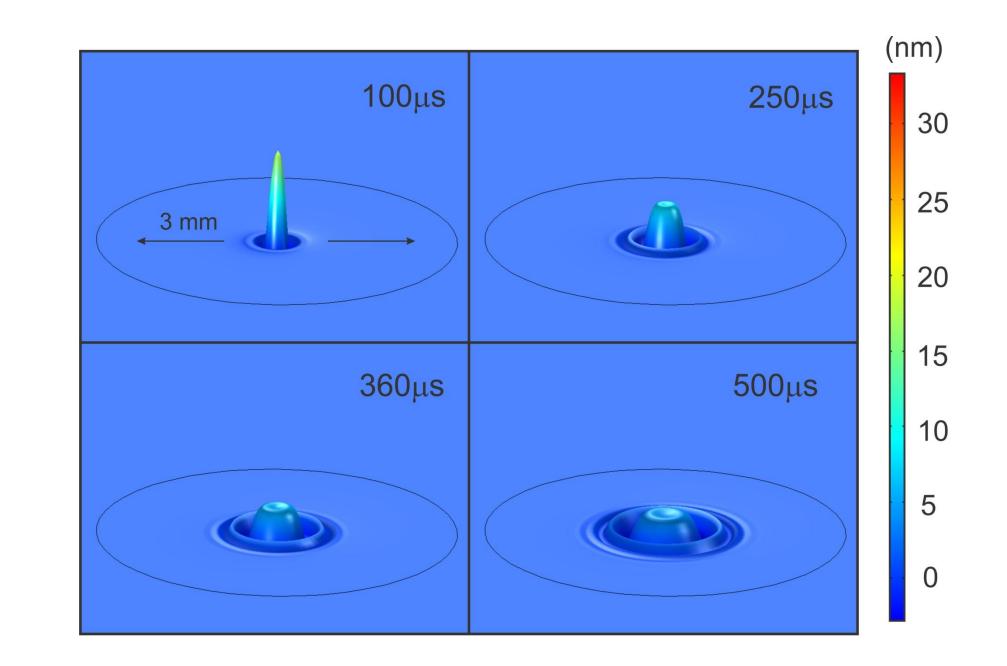


Figure 3. Water surface deformation under cw excitation.



Navier-Stokes : equation

$$\rho \frac{\partial \mathbf{u}}{\partial t} + \rho (\mathbf{u} \cdot \nabla) \mathbf{u} = -\nabla p + \mu \nabla^2 \mathbf{u} + \mathbf{F}$$
$$\nabla \cdot \mathbf{u} = 0$$

Volume force:

Pressure

radiation:
$$P_{cw}(r,t) = -\frac{2}{c} \left(\frac{n-1}{n+1}\right) \frac{2P_e}{\pi w_e^2} \exp\left(-\frac{2r^2}{w_e^2}\right)$$

 $(r,t) = -\frac{2}{c} \left(\frac{n-1}{n+1}\right) \frac{2Q}{\pi w_e^2} \exp\left(-\frac{2r^2}{w_e^2}\right) \frac{1}{t_0} \exp\left[-\frac{(t-\xi)^2}{\tau^2}\right]$

$$P_{Pulsed}(r,t) = -\frac{2}{c} \left(\frac{n-1}{n+1}\right) \frac{2Q}{\pi w_e^2} \exp\left(-\frac{2r^2}{w_e^2}\right) \frac{1}{t_0} \exp\left[-\frac{(t-\xi)^2}{\tau^2}\right]$$

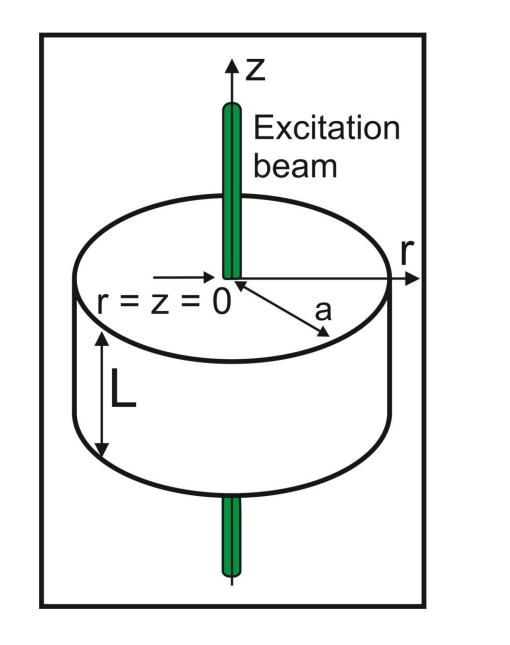


Figure 2. The model

was built in the 2D

axisymmetric

geometry.

Table 1. Parameters used for the simulations.

> Units Variable Value Kg/m³ 998.2 Density

Figure 4. Water surface deformation under pulsed excitation.

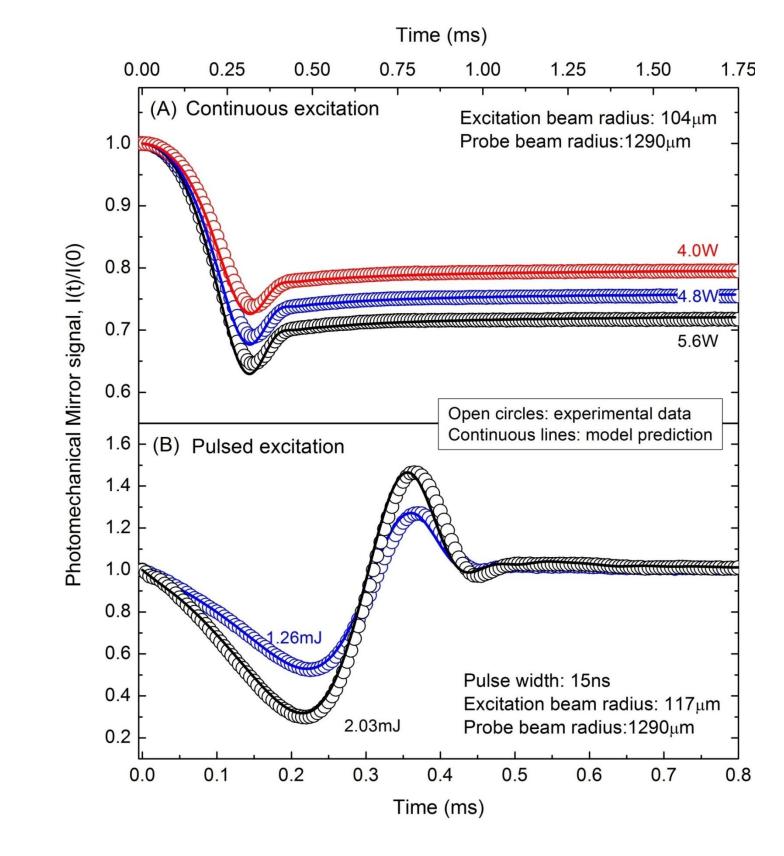


Figure 5. Time-resolved PM transients¹.

Conclusions: The numerical predictions are in

Dynamic 0.893 μ Viscosity Surface 72 σ Tension Refractive 1.33 n Index 30 Radius а

8 Thickness L mm

excellent agreement for both the continuous and pulsed excitation transients. In fact, it shows quantitatively that the effects of radiation forces in water can be fully described.

References:

- N. Astrath et al., Unraveling the effects of radiation forces in water, Nature Communications, 5, 5363 (2014).
- R. Pfeifer et al., Colloquium: Momentum of an 2. electromagnetic wave in dielectric media. Rev. Mod. Phys., 79, 1197-1216 (2007).

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10-3

Pa.s

N/m

mm